# The Crystal Structure of $\varepsilon$-Aminocaproic Acid* 

By Géza Bodor $\dagger$, Allan L. Bednowitz and Ben Post<br>Polytechnic Institute of Brooklyn, Brooklyn, New York, U.S.A. and Brookhaven National Laboratory, Upton, New York, U.S.A.

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The crystal structure of $\varepsilon$-aminocaproic acid, $\mathrm{NH}_{2}\left(\mathrm{CH}_{2}\right)_{5} \mathrm{COOH}$, which undergoes polycondensation in the solid state to form Nylon-6, has been determined. The unit cell is monoclinic with $a=8 \cdot 555$, $b=5.872, c=15.274 \AA ; \beta=103 \cdot 1^{\circ}$. The space group is $P 2_{1} / c$ and there are four molecules per unit cell. The signs of structure factors were determined by the symbolic addition procedure. Approximate atomic position coordinates were obtained from an $E$-factor map and the structure was refined by least-squares methods. The molecules are present in the crystal as zwitterions with three hydrogen atoms at $0.9 \AA$ from each nitrogen atom. Each zwitterion forms three $\mathrm{N}-\mathrm{H}-\mathrm{O}$ bonds with adjacent molecules to form a double-layered structure. The hydrogen bond lengths ( $\mathrm{N}-\mathrm{O}$ ) are 2.728, 2.753 and 2.814 , all $\pm 0.005 \AA$. Intramolecular bond lengths are: C-C 1.511 to $1.523 \AA, \mathrm{C}-\mathrm{O} 1.235$ and $1.261 \AA$, and $\mathrm{C}-\mathrm{N} 1.485 \AA$, all $\pm 0.004 \AA$. Bond angles are normal.

## Introduction

Muller (1952) has reported that Nylon-6, polypentamethylene capramide:

can be prepared by the solid state polycondensation of $\varepsilon$-aminocaproic acid. Recently Morosoff, Lim \& Morawetz (1964) showed that the polycondensation of single crystals of monomer yields highly oriented polycrystalline polymer specimens. Evidently the reactive functional groups in the monomer play important roles in determining the orientation as well as the nature of the reaction product. This investigation was undertaken to determine the crystal structure of the monomer and to study the relationships between the orientations of reactive groups in the monomer and in the polycondensation product, Nylon-6.

## Experimental

Crystals of $\varepsilon$-aminocaproic acid were obtained by evaporation from benzene solutions. Full data in three dimensions were collected with a counter diffractometer using filtered $\mathrm{Cu} K \alpha$ radiation. The crystal system is monoclinic; $a=8.855 \pm 0.002, b=5 \cdot 782 \pm 0.002, c=$ $15 \cdot 274 \pm 0 \cdot 004 \AA, \beta=103 \cdot 1 \pm \cdot 1^{\circ}$. The space group is $P 2_{1} / c$ and there are four molecules per unit cell. Of the 1485 independent reflections accessible to $\mathrm{Cu} K \alpha$ radiation, 1194 were measured and used in the structure determination. The others were either too weak

[^0]to permit reasonably precise measurement or were not detected at all and were given zero weight in the leastsquares analysis of the data. All 1194 reflection intensities were measured by the stationary-crystal station-ary-counter method, with a target take-off angle of $5^{\circ}$. Background measurements were made on high and low angle sides of the reflections and the averaged values were used to correct the intensities for background effects. In addition, approximately 150 reflections, randomly distributed throughout the range of Bragg angles investigated, were measured by the mov-ing-crystal moving-counter ( $\theta-2 \theta$ ) procedure. After being corrected for background effects, these 'integrated' intensities were plotted against the corresponding peak intensities as functions of the Bragg angle. The smoothed plot was used to convert all 'peak' intensities to 'integrated' intensities.

A trial scale factor and an isotropic molecular temperature factor were obtained by a statistical analysis of the intensity data in the manner described by Wilson (1942). The magnitudes of normalized structure factors $\left(\left|E_{\mathrm{H}}\right|\right)$ were then computed using the equation:

$$
\begin{equation*}
E_{\mathbf{H}}^{2}=F_{\mathbf{H}}^{2} \sum_{j=1}^{N} \varepsilon f_{j \mathbf{H}}^{2} \tag{1}
\end{equation*}
$$

where $N$ is the number of atoms in the unit cell and $\varepsilon$ equals 2.0 for $h 0 l$ and $0 k 0$ reflections and 1.0 for other reflections. Only normalized structure factors with magnitudes greater than $1 \cdot 2(20 \%$ of the total) were used in the phase analysis.

## Determination of the structure

The symbolic addition procedure (Karle \& Karle, 1963) was used to determine phases from the magnitudes of the normalized structure factors. The phase determining procedure was facilitated by the use of a computer program previously described (Bednowitz \& Post,

1966; Okaya \& Bednowitz, 1967). The principal phase relation used in this procedure is the $\Sigma_{2}$ formula of Hauptman \& Karle (1953):

$$
\begin{equation*}
s E_{\mathbf{H}} \approx s \sum E_{\mathbf{K}} E_{\mathbf{H}-\mathbf{K}} \tag{2}
\end{equation*}
$$

where $s$ means 'sign of'. The right hand side or the 'double product sum' is modified in our procedure so that the residue of the dominant symbol determines the correct symbol assignment. The dominant symbol is the one whose double product sum is greater than the sum of all other double products associated with all other symbols. The difference between the double product sum of the dominant symbol and the sum of the double product of all other symbols is the residue ( $R_{s}$ ). ( $R_{s}$ ) must ke less than or equal to the absolute magnitude of the double product sum when the actual signs ( + or - ) are used. By modifying the usual probability formula (Karle \& Karle, 1963; Woolfson, 1954) to

$$
\begin{equation*}
P_{s}\left(E_{\mathbf{H}}\right) \geq \frac{1}{2}+\frac{1}{2} \tanh \frac{\sigma_{3}}{\sigma_{2}^{3 / 2}}\left|E_{\mathbf{H}}\right| R_{s} \tag{3}
\end{equation*}
$$

automated decisions could be made on a reasonably quantitative basis.
In order to initiate the automatic phase determining procedure it is necessary to choose a set of starting phases. For space group $P 2_{1} / c$ these three reflections are given arbitrary signs to fix the origin. Additional reflections with relatively large $E$ factors, and with many interactions with other (strong) terms, are then assigned symbolic phases. The interaction list is usually computed before assigning the first set of phases. The starting set can be augmented by using additional phase relations. For example, for this structure the $\Sigma_{1}$ relation (4) (Hauptman \& Karle, 1953),

$$
\begin{equation*}
s\left(E_{2 \mathbf{H}}\right) \approx s\left(E_{\mathbf{H}}^{2}-1\right) \tag{4}
\end{equation*}
$$

was used to establish a positive sign for the reflection ( $46 \overline{8}$ ). The starting set used in the phase determination is listed in Table 1.

## Table 1. Starting set of phases



Using the starting set of six reflections and applying the symbolic addition procedure manually to the 68 reflections with $|E|>2 \cdot 0,11$ signs and 16 symbols were determined. In the next stage 108 reflections with $|E|>1.8$ were analyzed and yielded a cumulative total of 26 signs and 44 symbols. In order to facilitate the rapid development of more symbolic assignments, the symbol (c) was then assigned to (620).

In the second pass through the computer 199 reflections with $|E|>1.50$ were examined. A total of 42 signs and 123 symbols were known after this pass. For the third pass the minimum $|E|$ was reduced to 1.2 and 318 reflections were used and 49 signs and 193 symbols were established.

There were many indications that $(b=c)(a=-)$ and $(b=+)$ were the most probable sign assignments. Although additional phases could undoubtedly have been determined by substituting these signs for the symbols in the phase determining procedure and continuing that process, we decided to terminate the phase analysis and prepare $E$ maps based on signs and symbols determined at the end of the third computer pass. The nine nonhydrogen atoms in the molecule were clearly revealed

Table 2. Atomic coordinates and thermal parameters
Estimated standard errors of the coordinates, the $b_{i j}$ and their standard errors have been multiplied by $10^{4}$. The thermal parameters are given by: $h^{2} b_{11}+k^{2} b_{22}+l^{2} b_{33}+2 h k b_{12}+2 h l b_{13}+2 k l b_{23}$.

|  | $x$ | $\sigma_{x}$ | $y$ | $\sigma_{y}$ | $z$ | $\sigma_{z}$ | $B$ | $b_{11}$ | $b_{22}$ | $b_{33}$ | $b_{12}$ | $b_{13}$ | $b_{23}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{O}(1)$ | 0.6506 | 2 | 0.2975 | 2 | 0.0921 | 1 |  | 178 (3) | 126 (4) | 73 (1) | 43 (3) | 69 (1) | $17{ }^{(2)}$ |
| $\mathrm{O}(2)$ | $0 \cdot 5396$ | 2 | 0.0183 | 2 | 0.1531 | 1 |  | 123 (2) | 165 (4) | 41 (1) | -2 (2) | 31 (1) | -10 (1) |
| C(1) | 0.6327 | 2 | 0.0932 | 3 | $0 \cdot 1065$ | 1 |  | 79 (3) | 150 (6) | 33 (1) | 25 (3) | 10 (1) | -2 (2) |
| C(2) | 0.7333 | 2 | 0.9210 | 4 | 0.0688 | 1 |  | 93 (3) | 174 (6) | 40 (1) | 29 (3) | 17 (1) | -12(2) |
| C(3) | 0.9012 | 2 | 0.9039 | 4 | $0 \cdot 1285$ | 2 |  | 102 (3) | 263 (8) | 42 (1) | 61 (4) | 10 (1) | -21 (2) |
| C(4) | $0 \cdot 0047$ | 2 | 0.7282 | 4 | 0.0951 | 1 |  | 95 (2) | 195 (7) | 46 (1) | 37 (3) | 17 (1) | -14(2) |
| C(5) | $0 \cdot 1775$ | 2 | 0.7302 | 4 | $0 \cdot 1480$ | 1 |  | 103 (3) | 182 (6) | 39 (1) | 46 (3) | 18 (1) | -17(2) |
| C(6) | $0 \cdot 2791$ | 2 | 0.5532 | 4 | $0 \cdot 1152$ | 1 |  | 92 (3) | 157 (6) | 43 (1) | 24 (3) | 18 (1) | -10 (2) |
| N | $0 \cdot 4502$ | 2 | 0.5745 | 3 | $0 \cdot 1627$ | 1 |  | 99 (2) | 143 (5) | 42 (1) | 28 (3) | 24 (1) | 7 (2) |
| H(1) | 0.7438 | 26 | 0.9662 | 40 | 0.0093 | 16 | $3 \cdot 3$ (5) |  |  |  |  |  |  |
| H(2) | $0 \cdot 6808$ | 27 | 0.7729 | 45 | 0.0650 | 16 | $3 \cdot 4$ (4) |  |  |  |  |  |  |
| H(3) | 0.9639 | 29 | 0.0555 | 48 | 0.1367 | 16 | $4 \cdot 1$ (5) |  |  |  |  |  |  |
| H(4) | $0 \cdot 8944$ | 30 | 0.8790 | 50 | 0. 1900 | 20 | $4 \cdot 9$ (6) |  |  |  |  |  |  |
| H(5) | 0.0081 | 30 | 0.7667 | 50 | 0.0275 | 18 | $4 \cdot 6$ (6) |  |  |  |  |  |  |
| H(6) | 0.9617 | 24 | $0 \cdot 5975$ | 42 | $0 \cdot 1022$ | 14 | $2 \cdot 4$ (4) |  |  |  |  |  |  |
| H(7) | $0 \cdot 1836$ | 30 | 0.7120 | 46 | 0.2157 | 19 | $4 \cdot 7$ (6) |  |  |  |  |  |  |
| H(8) | $0 \cdot 2214$ | 30 | 0.8884 | 49 | $0 \cdot 1404$ | 17 | $4 \cdot 4$ (6) |  |  |  |  |  |  |
| H(9) | $0 \cdot 2414$ | 26 | 0.4149 | 42 | $0 \cdot 1207$ | 14 | $2 \cdot 8$ (5) |  |  |  |  |  |  |
| $\mathrm{H}(10)$ | $0 \cdot 2829$ | 27 | $0 \cdot 5877$ | 44 | 0.0518 | 17 | $3 \cdot 5$ (5) |  |  |  |  |  |  |
| $\mathrm{H}(11)$ | $0 \cdot 4551$ | 25 | $0 \cdot 5571$ | 40 | 0.2241 | 16 | $3 \cdot 2$ (5) |  |  |  |  |  |  |
| H(12) | 0.4939 | 26 | 0.7142 | 48 | $0 \cdot 1562$ | 15 | $3 \cdot 4$ (5) |  |  |  |  |  |  |
| H(13) | $0 \cdot 5181$ | 27 | $0 \cdot 4651$ | 42 | $0 \cdot 1361$ | 15 | $3 \cdot 3$ (5) |  |  |  |  |  |  |

Table 3. $\left|F_{\text {obs }}\right|$ and $\left|F_{\text {calc }}\right|$


Table 3 (cont.)







 22
33
76
13
24 22
33
26
13
24 $H=$
$K=$
0
33
72
53
141
16
118
42
55
45
0
47
56
36
90



AL
$L$ fabs fCAL
on one of the maps and the signs used to compute that map were used to refine the atomic coordinates. The indicated positions of these atoms were taken as the starting points of a least-squares analysis, using the set of 242 reflections for which phases had been established. Initially, the atoms were assumed to have isotropic thermal motion. The residual $(R)$,

$$
R=\frac{\sum| | F_{\mathrm{obs}}\left|-\left|K F_{\mathrm{catc}}\right|\right|}{\sum\left|F_{\mathrm{obs}}\right|}
$$

for this limited set of data was reduced to $0 \cdot 128$ before including additional reflection data in the least-squares analysis. Because of limitations in the least-squares program and in the programmed data storage area, the data were split arbitrarily into two groups, each containing about 750 reflections and refined separately to residuals of 0.22 and 0.20 , with only slight differences ( $0.001-0.0012$ ) among the atomic coordinates derived from the two sets separately.

Refinement was then continued with the use of all the data. The isotropic refinement converged to $R=$ $0 \cdot 15$. After the introduction of anisotropic temperature factors, $R$ decreased to $0 \cdot 12$. At this point an electron density difference map was computed. A composite drawing, based on this map, revealed the locations of all thirteen hydrogen atoms in the asymmetric unit of structure. It was also clear at this stage that one hydrogen atom had been shifted from the carboxyl to the amine group to form a zwitterion. After introduction of the hydrogen atoms into the least-squares analysis with fixed and finally with variable isotropic temperature factors, the residual was reduced to 0.05 (not including the unobserved reflections). The final atomic parameters and temperature factors are listed in Table 2. The observed and calculated structure factors are listed in Table 3; the unobserved reflections (given zero weight in the least-square analysis) are indicated by asterisks. The estimation of $\sigma$ used in the standard weighting scheme was:

$$
\begin{array}{ll}
F_{\text {obs }}>10 \cdot 0 & \sigma=0 \cdot 1 F_{\text {obs }} \\
F_{\text {obs }} \leq 10 \cdot 0 & \sigma=1 \cdot 0 \\
F_{\text {obs }}=0 \cdot 0 & \text { weight }=0 .
\end{array}
$$

## Discussion

## Accuracy of the structure

The refinement of the structure was continued until changes of coordinates of non-hydrogen atoms between successive least-squares cycles were less than two per cent of their standard errors; the corresponding changes for hydrogen atoms were less than four per cent.

Standard errors of parameters were computed by inversion of the least-squares matrices. The indicated values were then multiplied by two and rounded off to the nearest significant digit before being included in this manuscript. Hamilton (1965) has emphasized the usefulness of a procedure of this sort to furnish
a measure of the accuracy, rather than the precision, of a crystal structure determination. The final value of $R$ for the 1194 observed reflections was $5 \cdot 0 \%$.

## Temperature factors

Temperature factors are listed in Tables 2 and 6. The magnitudes of atomic thermal motions along the principal axes of the thermal ellipsoids are given in Table 6 together with the r.m.s. thermal displacements of the hydrogen atoms. As expected, the thermal motions of the carbon atoms are smallest along the molecular chain direction. The magnitudes of the thermal motions of the atoms in the molecular chain are largest near the center of the chain, i.e. at $C(3)$ and $C(4)$. The hydrogen bonds at both ends of the molecules restrict the freedom of movement of atoms nearest those bonds and have least effect on atoms near the chain centers. Only axis 3 of $\mathrm{O}(1)$ does not fit this description; its anomalously large value may reflect the fact that it is associated with only one hydrogen bond; $\mathrm{O}(2)$ is involved in two such bonds.

## Molecular configuration

Fig. 1 is a schematic drawing of the molecule, listing bond lengths and approximate magnitudes and orientations of the principal axes of the atomic thermal ellipsoids relative to the long molecular direction. Intramolecular bond lengths and intermolecular $\mathrm{N}-\mathrm{H}-\mathrm{O}$ distances are listed in Tables 4 and 5. As might have been expected, the two carbon-oxygen bonds in the molecule are almost equal in length, 1.235 and $1.261 \AA$ both $\pm 0.005 \AA$; after correction for 'riding' of the oxygen atoms on the carbon atoms to which they are bonded (Busing \& Levy, 1964) these become 1.263 and $1 \cdot 271 \AA$. Similar results have been reported by Marsh and co-workers in their very precise analyses of the zwitterions l-lysine monochloride dihydrate (Wright \& Marsh, 1962) with C-O distances of 1.246 and $1 \cdot 250 \AA$ and L-alanine (Simpson \& Marsh, 1966) with 1.253 and 1.265 Å.

Table 4. Bond lengths and angles

| Bond lengths* |  | Bond angles $\ddagger$ |  |
| :--- | :--- | :--- | :--- |
| $\mathrm{C}(1)-\mathrm{O}(1)$ | $1.235 \AA \AA \dagger$ | $\mathrm{O}(1)-\mathrm{C}(1)-\mathrm{O}(2)$ | $124 \cdot 0^{\circ}$ |
| $\mathrm{C}(1)-\mathrm{O}(2)$ | $1.261 \dagger$ | $\mathrm{O}(1)-\mathrm{C}(1)-\mathrm{C}(2)$ | 118.3 |
| $\mathrm{C}(1)-\mathrm{C}(2)$ | 1.523 | $\mathrm{O}(2)-\mathrm{C}(1)-\mathrm{C}(2)$ | 17.7 |
| $\mathrm{C}(2)-\mathrm{C}(3)$ | 1.520 | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{C}(3)$ | $110 \cdot 8$ |
| $\mathrm{C}(3)-\mathrm{C}(4)$ | 1.526 | $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{C}(4)$ | 112.7 |
| $\mathrm{C}(4)-\mathrm{C}(5)$ | 1.516 | $\mathrm{C}(3)-\mathrm{C}(4)-\mathrm{C}(5)$ | 112.8 |
| $\mathrm{C}(5)-\mathrm{C}(6)$ | 1.511 | $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{C}(5)$ | 12.5 |
| (Average C-C C | $1.519)$ | $\mathrm{C}(5)-\mathrm{C}(6)-\mathrm{N}$ | 110.9 |
|  |  | (Average C-C-C angle |  |
| $\mathrm{C}(6)-\mathrm{N}$ | 1.485 |  | $\left.112.2^{\circ}\right)$ |
| N-H-O(1) | 2.753 |  |  |
| N-H-O(2) | 2.728 |  |  |
| N-H-O(2') | 2.814 |  |  |

[^1]The average length of the five carbon-carbon bonds in $\varepsilon$-aminocaproic acid is $1.519 \AA$. For the two crystal structures mentioned above, Marsh and co-workers reported average carbon-carbon bond lengths of 1.523 and $1.525 \AA$. It appears that the 'normal' single carbon-
carbon bond length, at least for molecules of this type, is in the range 1.519 to $1.525 \AA$.

Hahn (1957) reported an apparent lengthening of the carbon-nitrogen $\left(\mathrm{C}-\mathrm{NH}_{3}^{+}\right)$bond in a number of zwitterions; he listed an average value of $1.503 \AA$ com-

Table 5. Bond lengths and angles involving hydrogen atoms

| Lengths* | Angles $\dagger$ |  | Angles $\dagger$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{C}(2)-\mathrm{H}(1) \quad 0.97 \AA$ | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{H}(1)$ | $109^{\circ}$ | $\mathrm{N}-\mathrm{C}(6)-\mathrm{H}(9)$ | $112^{\circ}$ |
| $\mathrm{C}(2)-\mathrm{H}(2) \quad 0.97$ | $\mathrm{C}(3)-\mathrm{C}(2)-\mathrm{H}(1)$ | 108 | $\mathrm{N}-\mathrm{C}-\mathrm{H}(10)$ | 102 |
| $\mathrm{C}(3)-\mathrm{H}(3) \quad 1.03$ | $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{H}(2)$ | 111 |  |  |
| $\mathrm{C}(3)-\mathrm{H}(4) \quad 0.97$ | $\mathrm{C}(3)-\mathrm{C}(2)-\mathrm{H}(2)$ | 111 | $\mathrm{H}(12)-\mathrm{N}-\mathrm{H}(11)$ | 106 |
| $\mathrm{C}(4)-\mathrm{H}(5) \quad 1.07$ | $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{H}(3)$ | 110 | $\mathrm{H}(13)-\mathrm{N}-\mathrm{H}(12)$ | 104 |
| $\mathrm{C}(4)-\mathrm{H}(6) \quad 0.92$ | $\mathrm{C}(4)-\mathrm{C}(3)-\mathrm{H}(3)$ | 107 | $\mathrm{H}(11)-\mathrm{N}-\mathrm{H}(13)$ | 116 |
| $\mathrm{C}(5)-\mathrm{H}(7) \quad 1.03$ | $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{H}(4)$ | 114 |  |  |
| $\mathrm{C}(5)-\mathrm{H}(8) \quad 1.02$ | $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{H}(4)$ | 113 | $\mathrm{H}(1)-\mathrm{C}(2)-\mathrm{H}(2)$ | 109 |
| $\mathrm{C}(6)-\mathrm{H}(9) \quad 0.89$ | $\mathrm{C}(3)-\mathrm{C}(4)-\mathrm{H}(5)$ | 109 | $\mathrm{H}(3)-\mathrm{C}(3)-\mathrm{H}(4)$ | 99 |
| $\mathrm{C}(6)-\mathrm{H}(10) \quad 1.00$ | $\mathrm{C}(5)-\mathrm{C}(4)-\mathrm{H}(5)$ | 106 | $\mathrm{H}(5)-\mathrm{C}(4)-\mathrm{H}(6)$ | 114 |
|  | $\mathrm{C}(3)-\mathrm{C}(4)-\mathrm{H}(6)$ | 107 | $\mathrm{H}(7)-\mathrm{C}(5)-\mathrm{H}(8)$ | 106 |
| (Average $\mathrm{C}-\mathrm{H}=0.99 \AA$ ) | $\mathrm{C}(5)-\mathrm{C}(4)-\mathrm{H}(6)$ | 108 | $\mathrm{H}(9)-\mathrm{C}(6)-\mathrm{H}(10)$ | 112 |
| $\mathrm{N}-\mathrm{H}(11) \quad 0.94$ | $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{H}(7)$ | 111 |  |  |
| $\mathrm{N}-\mathrm{H}(12) \quad 0.92$ | $\mathrm{C}(6)-\mathrm{C}(5)-\mathrm{H}(7)$ | 111 | (Average $\mathrm{H}-\mathrm{C}-\mathrm{H}$ | $108^{\circ}$ ) |
| $\mathrm{N}-\mathrm{H}(13) \quad 1.01$ | $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{H}(8)$ | 107 |  |  |
| (Average $\mathrm{N}-\mathrm{H}$ | $\mathrm{C}(6)-\mathrm{C}(5)-\mathrm{H}(8)$ | 110 | (hydrogen bonds) $\mathrm{N}-\mathrm{H}(13)-\mathrm{O}(1)$ | 174 |
|  | $\mathrm{C}(5)-\mathrm{C}(6)-\mathrm{H}(9)$ | 110 | $\mathrm{N}-\mathrm{H}(12)-\mathrm{O}(2)$ | 166 |
|  | $\mathrm{C}(5)-\mathrm{C}(6)-\mathrm{H}(10)$ | 109 | $\mathrm{N}-\mathrm{H}(11)-\mathrm{O}\left(\frac{1}{2}\right)$ | 179 |
|  | (Average $\mathrm{C}-\mathrm{C}-\mathrm{H}$ | 109.5 ${ }^{\circ}$ ) |  |  |

* The standard deviations of the bond lengths are $0.05 \AA$, except $0.06 \AA$ for $\mathrm{C}(3)-\mathrm{H}(3)$ and $\mathrm{C}(3)-\mathrm{H}(4)$.
$\dagger$ The standard deviation of the bond angles involving one hydrogen atom is $3^{\circ}$; for those involving two hydrogen atoms it is $4^{\circ}$.
Table 6. Thermal motions and orientations of thermal ellipsoids
R.M.S. components of thermal motion along principal axes ( $\AA$ )

| R.M.S. components of thermal motion along principal axes ( $\AA$ ) |  |  |  | Angles between principal axes and ve defined by $\mathrm{C}(2)-\mathrm{C}(6)\left(^{\circ}\right)$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Axis 1 | Axis 2 | Axis 3 | $\sqrt{\overline{\mu^{2}}}(\AA)$ | Axis 1 | Axis 2 | Axis 3 |
| 0.138 | $0 \cdot 194$ | 0.312 | 0.39 | 132 | 56 | 61 |
| 0.164 | $0 \cdot 184$ | $0 \cdot 229$ | 0.34 | 127 | 64 | 49 |
| 0.143 | $0 \cdot 181$ | $0 \cdot 196$ | $0 \cdot 30$ | 24 | 114 | 94 |
| $0 \cdot 148$ | 0.196 | 0.217 | $0 \cdot 33$ | 30 | 118 | 82 |
| $0 \cdot 151$ | $0 \cdot 204$ | 0.256 | 0.36 | 21 | 110 | 95 |
| 0.149 | $0 \cdot 205$ | 0.234 | $0 \cdot 34$ | 24 | 114 | 85 |
| 0.149 | $0 \cdot 210$ | 0.213 | 0.33 | 26 | 92 | 64 |
| 0.146 | $0 \cdot 191$ | $0 \cdot 221$ | 0.33 | 32 | 120 | 80 |
| $0 \cdot 146$ | $0 \cdot 184$ | $0 \cdot 221$ | $0 \cdot 32$ | 144 | 61 | 71 |



Fig. 1. Schematic drawing of the $\varepsilon$-aminocaproic acid molecule.

EPSILON-AMINOCAPROIC ACID
Fig. 2. Stereo pair view along the $c$ axis.
pared with a normal carbon-nitrogen single bond length of $1.47 \AA$. In $\varepsilon$-aminocaproic acid this distance is $1.485 \pm 0.005 \AA$; in the structures mentioned above, Marsh and co-workers reported corresponding values of $1.480,1.484$ and $1.491 \AA$. In view of the relatively high accuracy of these later values, it seems likely that the lengthening of this bond above the normal value is a real effect but one which is substantially smaller than indicated by the values quoted by Hahn.

The three hydrogen bonds formed by each nitrogen atom are: $2.753[\mathrm{~N}-\mathrm{O}(1)], 2.728[\mathrm{~N}-\mathrm{O}(2)]$ and $2.814 \AA$ $\left[\mathrm{N}-\mathrm{O}\left(2^{\prime}\right)\right]$, all $\pm 0.005 \AA$. It is evident that the fact that $\mathrm{O}(2)$ participates in two hydrogen bonds, while $\mathrm{O}(1)$ is involved in only one, does not affect the lengths of the hydrogen bonds significantly.

The average distance between hydrogen atoms bonded to the same carbon atom is $1.59 \pm 0.065 \AA$; the corresponding average distance among the hydrogen atoms in the $\mathrm{NH}_{3}^{+}$group is $1.55 \pm 0.065 \AA$.

Bond angles in the molecule are 'normal', i.e. similar to those reported for molecules of this general type. The average $\mathrm{C}-\mathrm{C}-\mathrm{C}$ angle is $112 \cdot 2^{\circ}$.

The spatial arrangement of the non-hydrogen atoms is conveniently described in terms of the dihedral angles between the planes of adjacent triatomic groups. Thus, the angle between $\mathrm{O}(1), \mathrm{O}(2), \mathrm{C}(1)$ and $\mathrm{C}(1), \mathrm{C}(2)$, $\mathrm{C}(3)$ is $98.4^{\circ}$; it is $7.6^{\circ}$ between $\mathrm{C}(1), \mathrm{C}(2), \mathrm{C}(3)$ and $C(3), C(4), C(5)$ and $5.6^{\circ}$ between the latter group and $\mathrm{C}(5), \mathrm{C}(6), \mathrm{N} \ldots$ all $\pm 0 \cdot 2^{\circ}$.


Fig. 3. Molecular layer in Nylon-6.

## Molecular packing

A stereo pair drawing (Johnson, 1965) of a portion of the unit cell, viewed along $\mathbf{c}$, is shown in Fig. 2. Two sets of molecular chains may be distinguished; one slopes down to the right along [110] and is inclined approximately $27 \frac{1}{2}^{\circ}$ to the horizontal axis a; the other is oriented along [ $1 \overline{1} 0]$. Each set lies in one layer. There are four such layers in each unit cell, all roughly perpendicular to $\mathbf{c}$.

Each nitrogen atom forms three hydrogen bonds; two link it with oxygen atoms $\left[\mathrm{O}(1)\right.$ and $\left.\mathrm{O}\left(2^{\prime}\right)\right]$ in molecules in the layer containing the nitrogen atom. The third bond is directed to an oxygen atom $\mathrm{O}\left(2^{\prime \prime}\right)$ in an adjacent layer.
Hydrogen bonds of the latter type link the layer at $z=\frac{1}{8}$ to that at $\frac{3}{8}$ and the layer at $z=\frac{5}{8}$ to that at $\frac{7}{8}$, forming strongly bonded double layers. These in turn are held together by relatively weak intermolecular interactions.

## The solid state conversion to Nylon-6

We have mentioned above that $\varepsilon$-aminocaproic acid undergoes polycondensation in the solid state. Under the combined influence of heat and high vacuum single crystals of the monomer transform to highly crystalline, oriented Nylon-6. It is possible to relate some structural features of the latter to the crystal structure of the monomer.

The crystal structure of Nylon-6 has been determined by Holmes, Bunn \& Smith (1955). The unit cell is monoclinic with $a=9.56, b=17 \cdot 2, c=8.01$ and $\beta=$ $112.5^{\circ}$; the space group is $P 2_{1}$ and there are eight 'condensed' monomer units in the unit cell. The structure may be described in terms of polymeric chains linked to one another by $\mathrm{N}-\mathrm{H}-\mathrm{O}$ bonds to form layers parallel to $b c$. The chains are parallel to $\mathbf{b}$ (Fig.3).

Examination of Fig. 3 reveals that the $\mathrm{C}=\mathrm{O}, \mathrm{N}, \mathrm{C}$ sequence in one chain is reversed in successive chains within the layer. It will be recalled that in the monomer the reactive groups (i.e. carboxyl and $\mathrm{NH}_{3}^{+}$) and the carbon atoms are arranged in the same sequence in all molecules in a given layer. The sequence is reversed in adjacent layers. It is evident that individual polymeric chains of Fig. 3 are derived from molecules lying within one layer in the crystal. At some stage the hydrogen bonds are broken and water is split out, leaving one hydrogen atom bonded to each nitrogen, and one oxygen atom in a carbonyl group. The carbon atom to which the oxygen is bonded $[\mathrm{C}(1)]$ forms a single bond link with the nitrogen atom.

Adjacent chains within one layer in Fig. 3 are derived from adjacent layers of the monomeric crystal. We have noted that the latter not only differ from one another in atomic sequence, but are inclined to each other by $55^{\circ}$. In the course of the polycondensation process these chains must change their relative orientations to become parallel to one another and to link up with one another by hydrogen bonding as shown in Fig. 3 .

Recent electron microscope studies of crystalline Nylon-6 specimens prepared from single crystals of the monomer reveal that twinning is of frequent occurrence in the polymer crystals (Fischer, 1965). The twinning appears to involve rotation of successive layers, or groups of layers, of the type shown in Fig.3, by $55^{\circ}$ relative to one another. Such twinning is not observed in Nylon specimens prepared from the melts (Macchi, 1966). Although the detailed mechanism of the twinning is not known, it is not difficult to relate the $55^{\circ}$ twin angle of the polymer to the $55^{\circ}$ by which molecular chains in successive layers of the monomer are inclined to one another.

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# The Crystal Structure of Thiodiglycollic Acid* 

By Sukla Paul $\ddagger \ddagger$<br>Indian Association for the Cultivation of Science, Jadavpur, Calcutta-32, India

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Thiodiglycollic acid, $\mathrm{S}\left(\mathrm{CH}_{2} \mathrm{COOH}\right)_{2}$, crystallizes in space group Pnam with 4 molecules in the unit cell of dimensions $a=5.050 \pm 0.005, b=6.701 \pm 0.003, c=17.742 \pm 0.004 \AA$. The structure has been solved from two-dimensional photographic data and has been refined by least squares with three-dimensional diffractometer data. The structure consists of parallel infinite chains of hydrogen bonded molecules. Hydrogen bonding takes place between carboxyl groups around the centre of symmetry; they are coplanar within experimental error.

## Introduction

Thiodiglycollic acid, $\mathrm{S}\left(\mathrm{CH}_{2} \mathrm{COOH}\right)_{2}$ (also known as methyl sulphide $\alpha, \alpha$-dicarboxylic acid or thiodiacetic acid), is of importance because of its use as a precipitating organic reagent for zirconium (Sant \& Sant, 1959). It is used for the detection of copper, lead, mercury and silver, and also for the estimation of cadmium.

A preliminary report of the structure has already been published (Roy, 1965).

## Experimental

Single crystals of thiodiglycollic acid were grown by slow evaporation of a concentrated aqueous solution

[^2]at room temperature. The crystals appeared as thin plates elongated in the [100] direction. The unit-cell dimensions determined by means of a General Electric Single Crystal Orienter are: $a=5 \cdot 050 \pm 0 \cdot 005, b=6 \cdot 701$ $\pm 0.003, c=17.742 \pm 0.004 \AA$.
Systematic absences as found on Weissenberg photographs about the $a$ and $b$ axes were consistent with the space groups Pnam and Pna2 ${ }_{1}$. The density as determined by flotation is $1.66 \mathrm{g.cm}^{-3}$, and the calculated density for four molecules per unit cell is $1.70 \mathrm{~g} . \mathrm{cm}^{-3}$.
The structure was solved from the two projections along [100] and [010]. The $0 k l$ and $h 0 l$ intensities were estimated visually from the integrated Weissenberg photographs. The three-dimensional data used for the refinement of the structure were collected with a General Electric diffractometer equipped with a scintillation counter using $\beta$-filtered $\mathrm{Cu} K \alpha$ radiation and discrimination.
Peak intensities were measured which were converted to integrated intensities in the following way. For a set of 30 reflexions of $\theta$ values, distributed over the accessible $\theta$ range, both peak intensities and inte-


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    $\dagger$ Present address: Research Institute for the Plastics Industry, Budapest, Hungary.

[^1]:    * The average standard deviation of the bond lengths is $0.005 \AA$; the largest is $0.006 \AA$, the smallest $0.004 \AA$.
    $\dagger$ After correction for 'riding' of the oxygen on $\mathrm{C}(1)$ these become $\mathrm{C}(1)-\mathrm{O}(1)=1 \cdot 263, \mathrm{C}(1)-\mathrm{O}(2)=1 \cdot 271 \AA$.
    $\ddagger$ The standard deviations of these bond angles range from 0.3 to $0.4^{\circ}$.

[^2]:    * Part of this work was done in the laboratory of the Department of Molecular Biophysics, Yale University, New Haven, Connecticut, U.S.A.
    $\dagger$ née Roy.
    $\ddagger$ Present address: Laboratory for Crystallography, University of Amsterdam, The Netherlands.

